

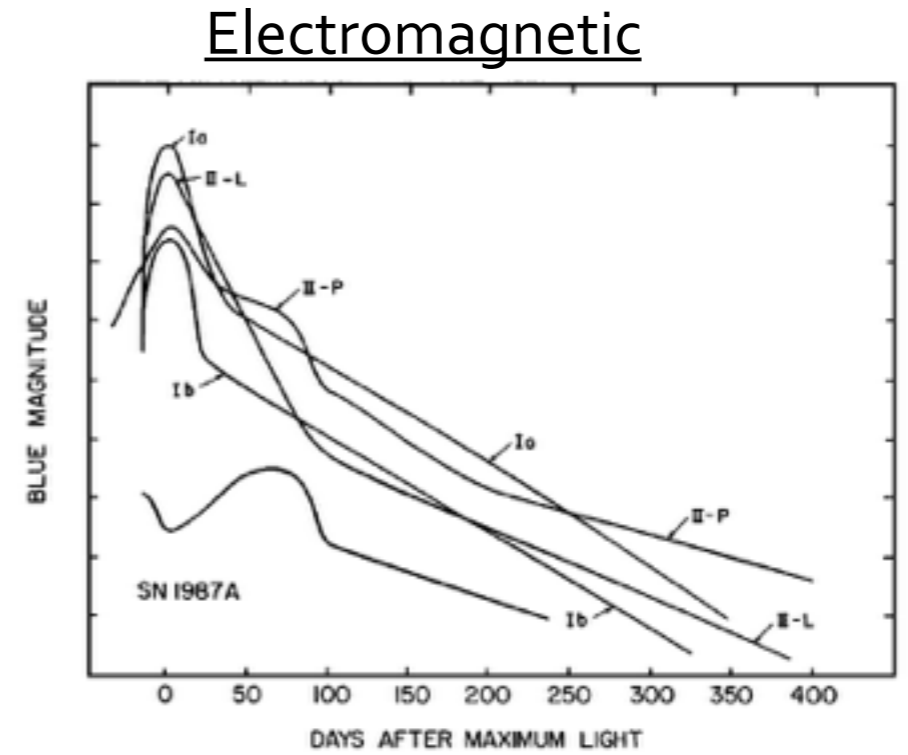
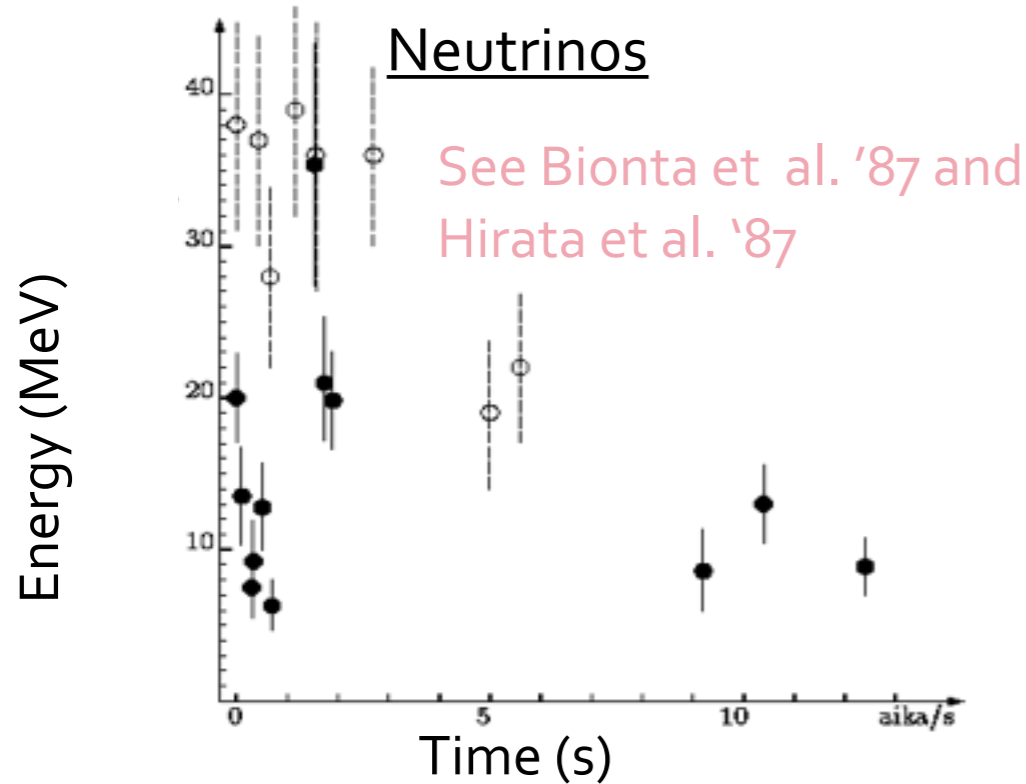
Three Dimensional Radiation Hydrodynamics Simulations of Core Collapse Supernovae

Luke F. Roberts

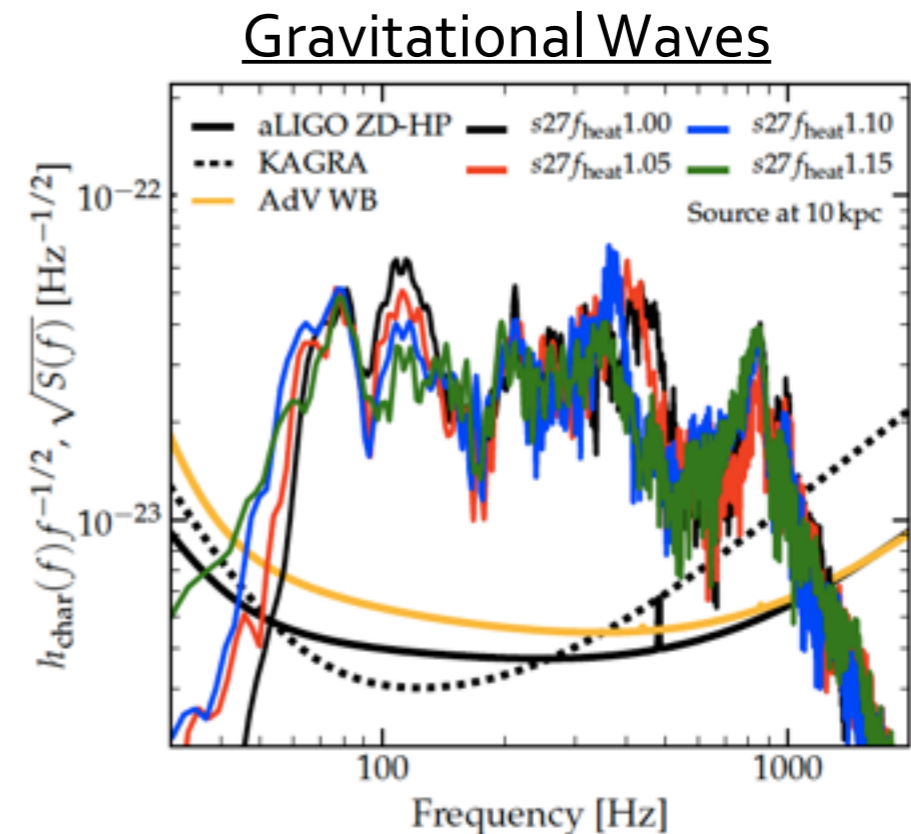
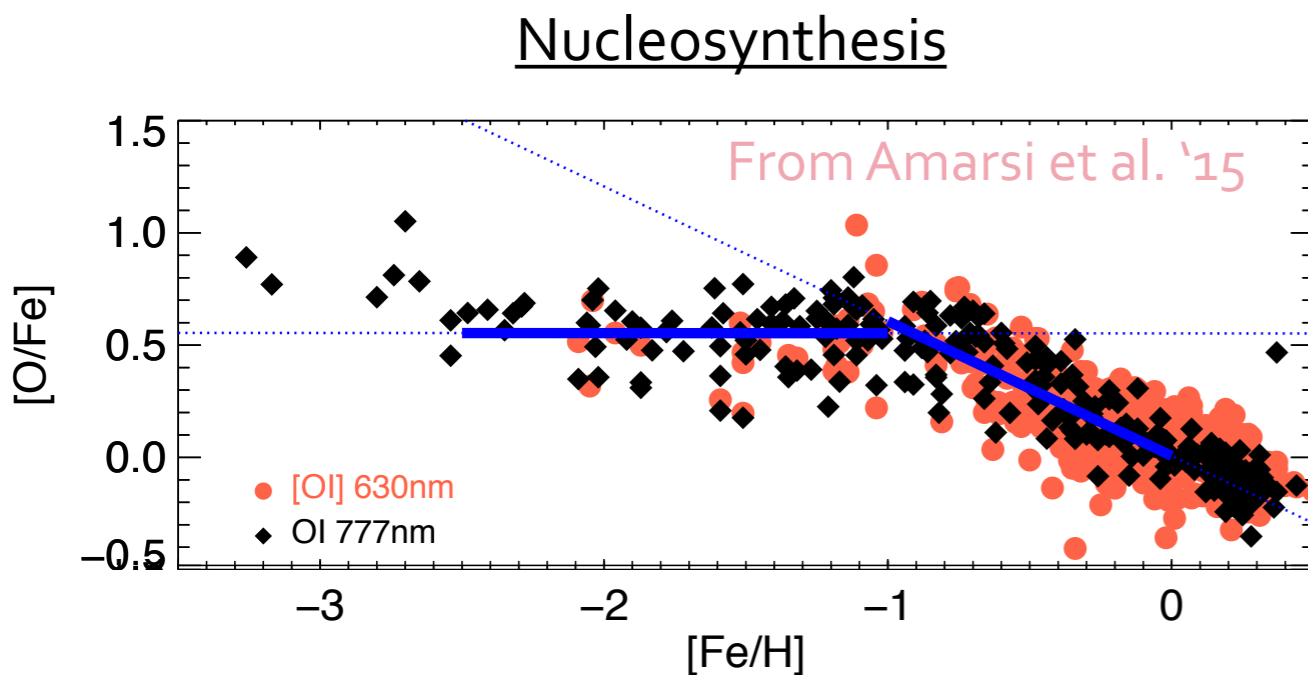
TAPIR, Caltech



Core Collapse Supernovae: Multi-messenger events



From Filippenko '97



From Ott et al. '12

Core Collapse

- Stars with $M > \sim 9 M_{\text{sun}}$ burn their core to Fe
- Core exceeds a Chandrasekhar mass \rightarrow supersonic collapse outside of homologous core \rightarrow bounce shock after $\sim 2 \times$ saturation density

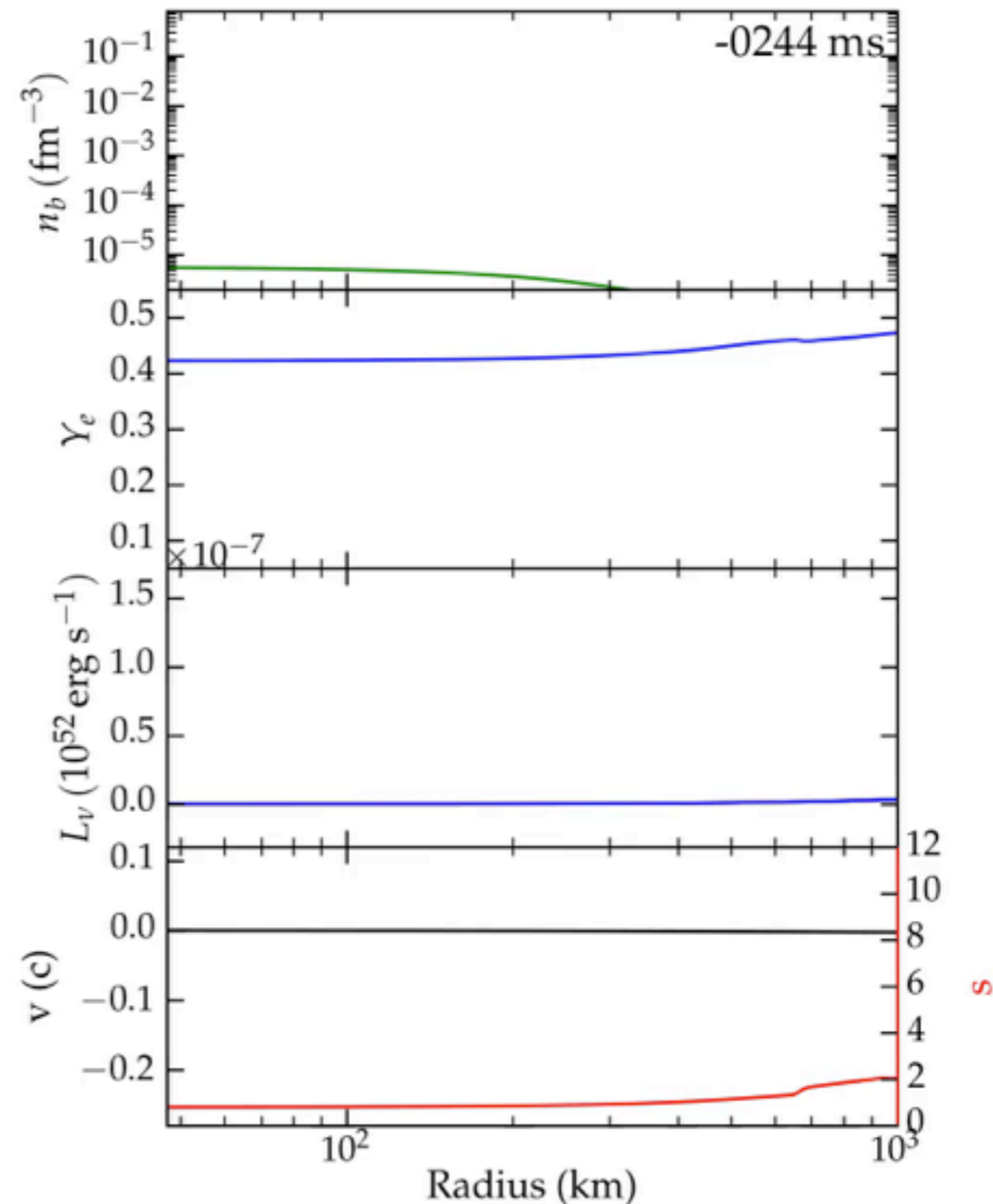
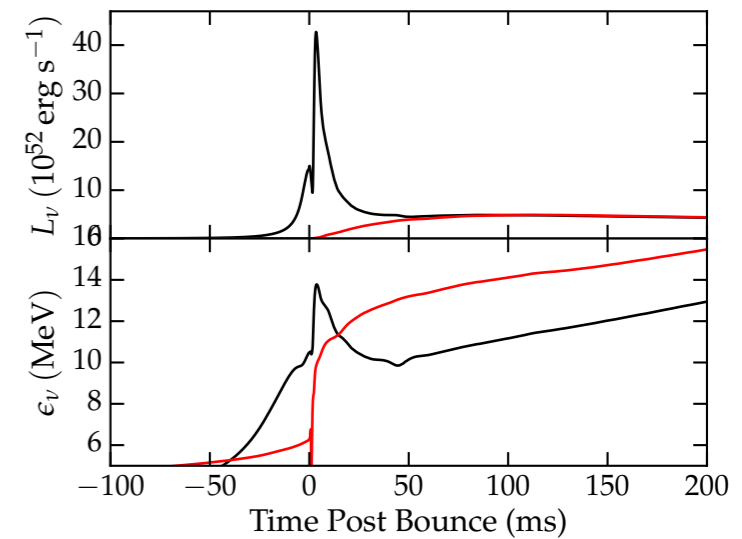
- Gravitational binding energy of compact remnant:

$$\frac{GM_{NS}^2}{R_{NS}} \sim 3 \times 10^{53} \text{ erg}$$

- Binding energy of stellar envelope:

$$\sim 10^{51} \text{ erg}$$

- No Explosions by the neutrino mechanism in spherical symmetry (e.g. Liebendorfer '00)



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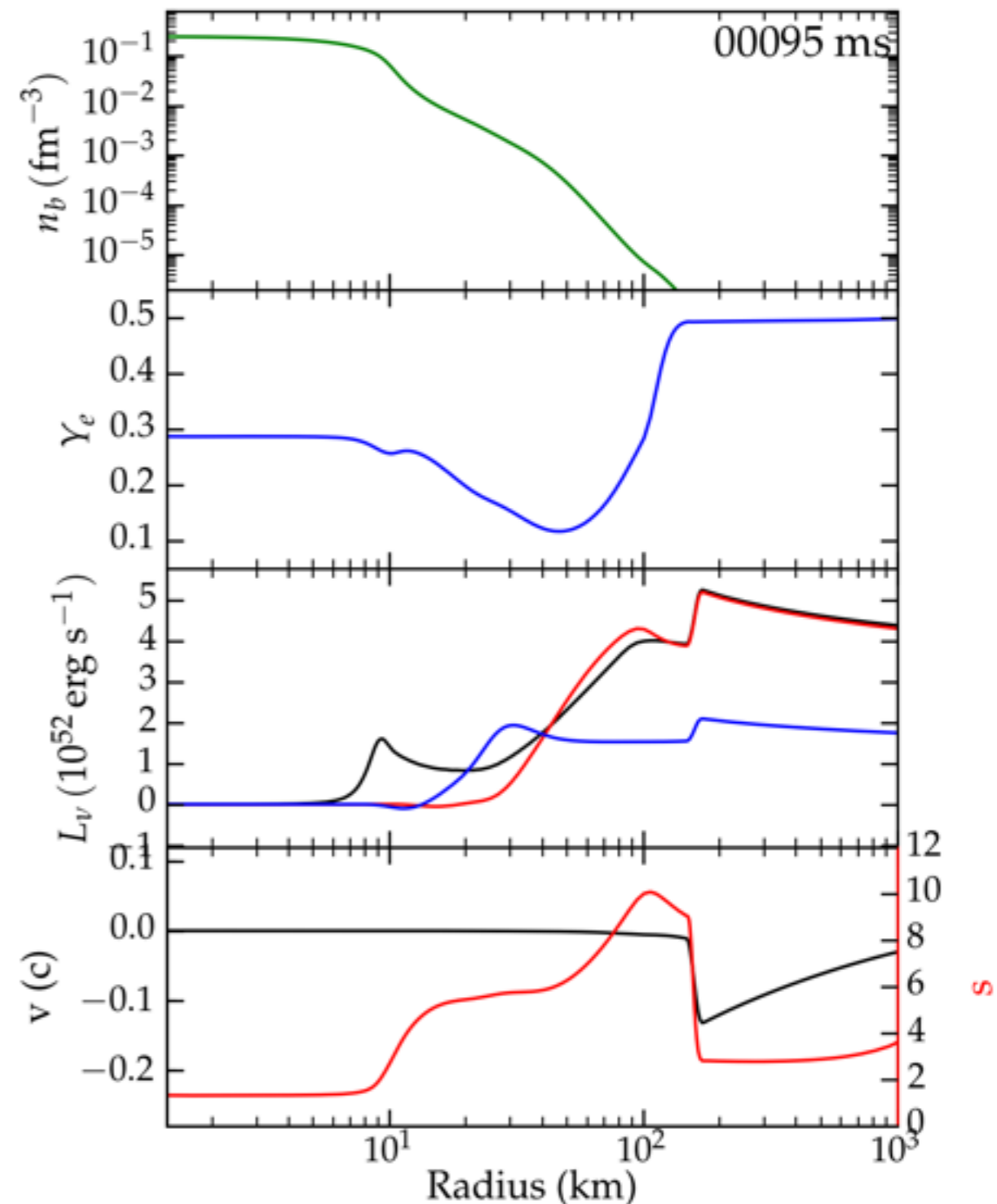
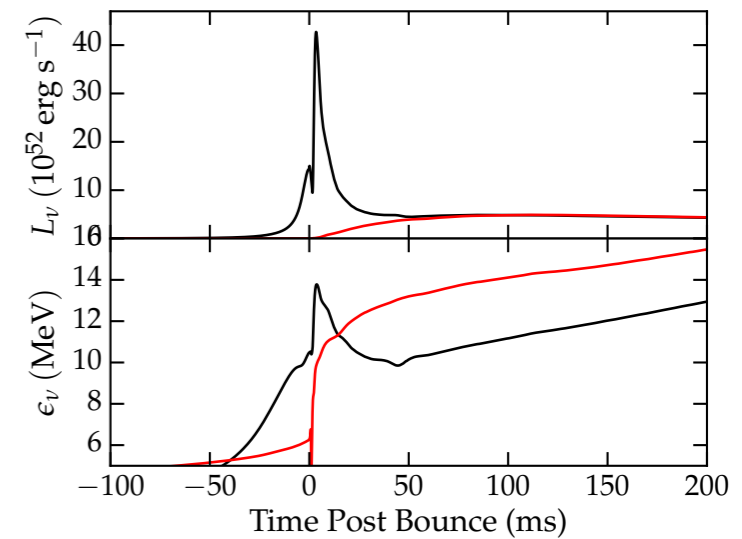
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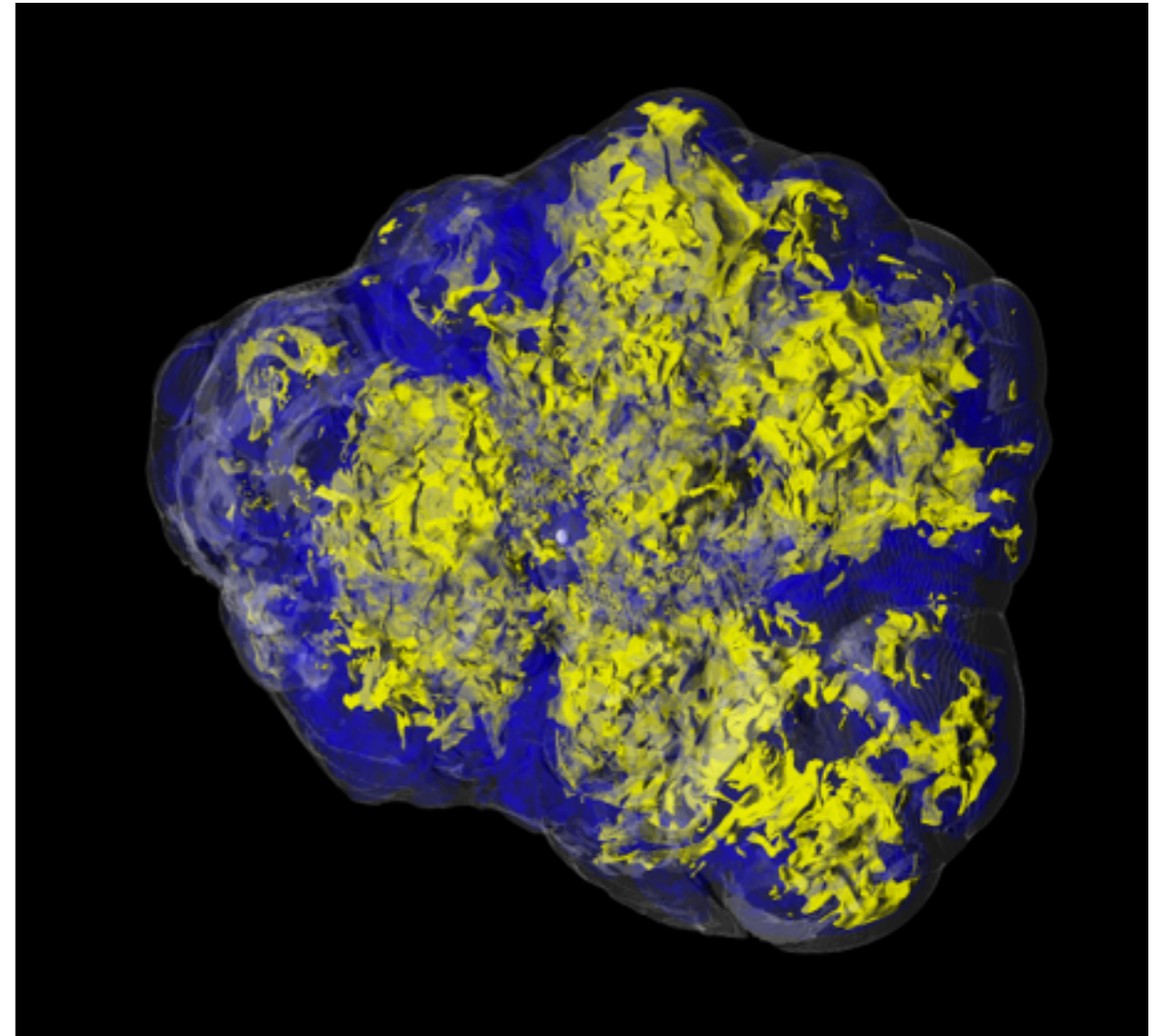
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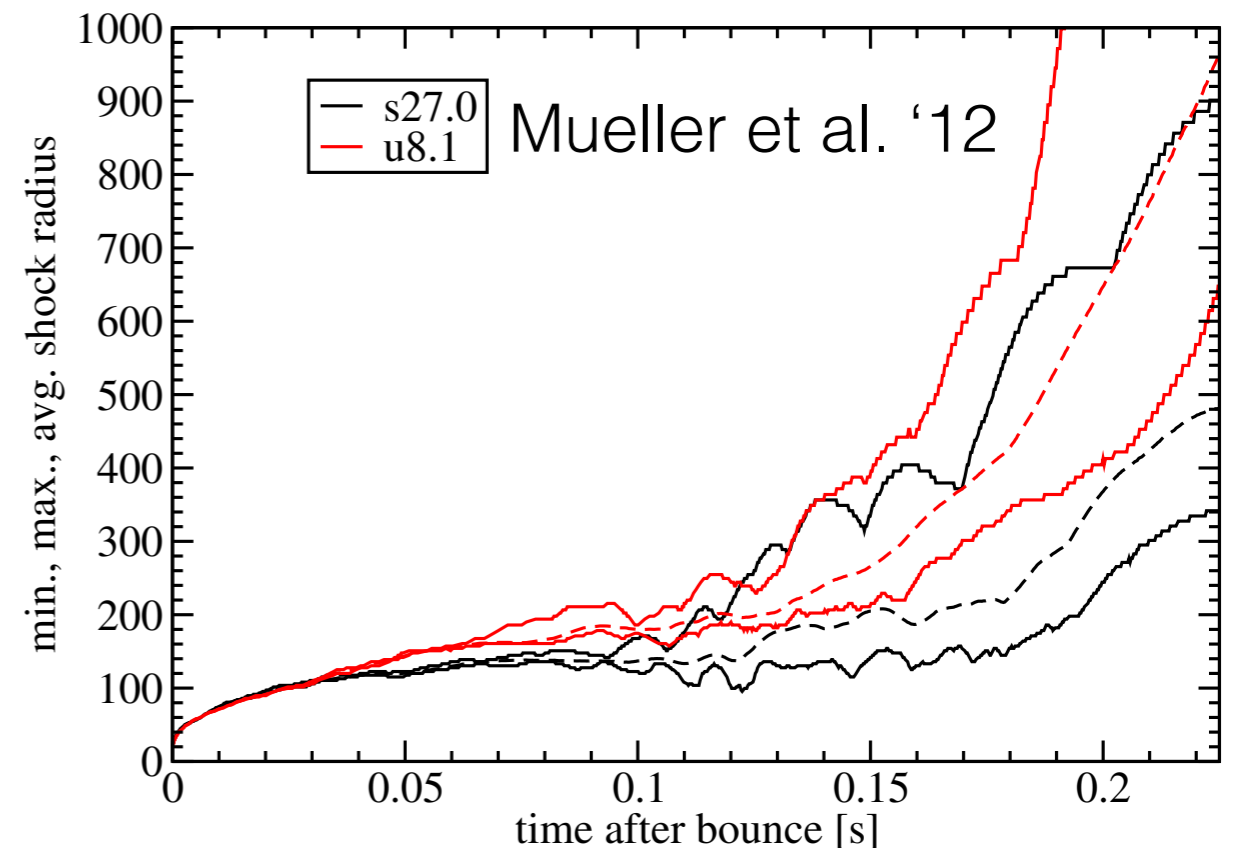
Post Bounce Evolution of CCSNe

- Hydrodynamic instabilities in the region behind the shock can provide increased post shock pressure and transport extra energy
- In axial symmetry, this enhances the efficiency of neutrino energy deposition and results in successful explosions (Mueller et al. '12, Bruenn et al. '13)
- Does the neutrino mechanism work in 3D? Lentz et al. '15 and Melson et al. '15 find it works for some progenitor stars using ray-by-ray transport
- How does this depend on input physics and numerics?



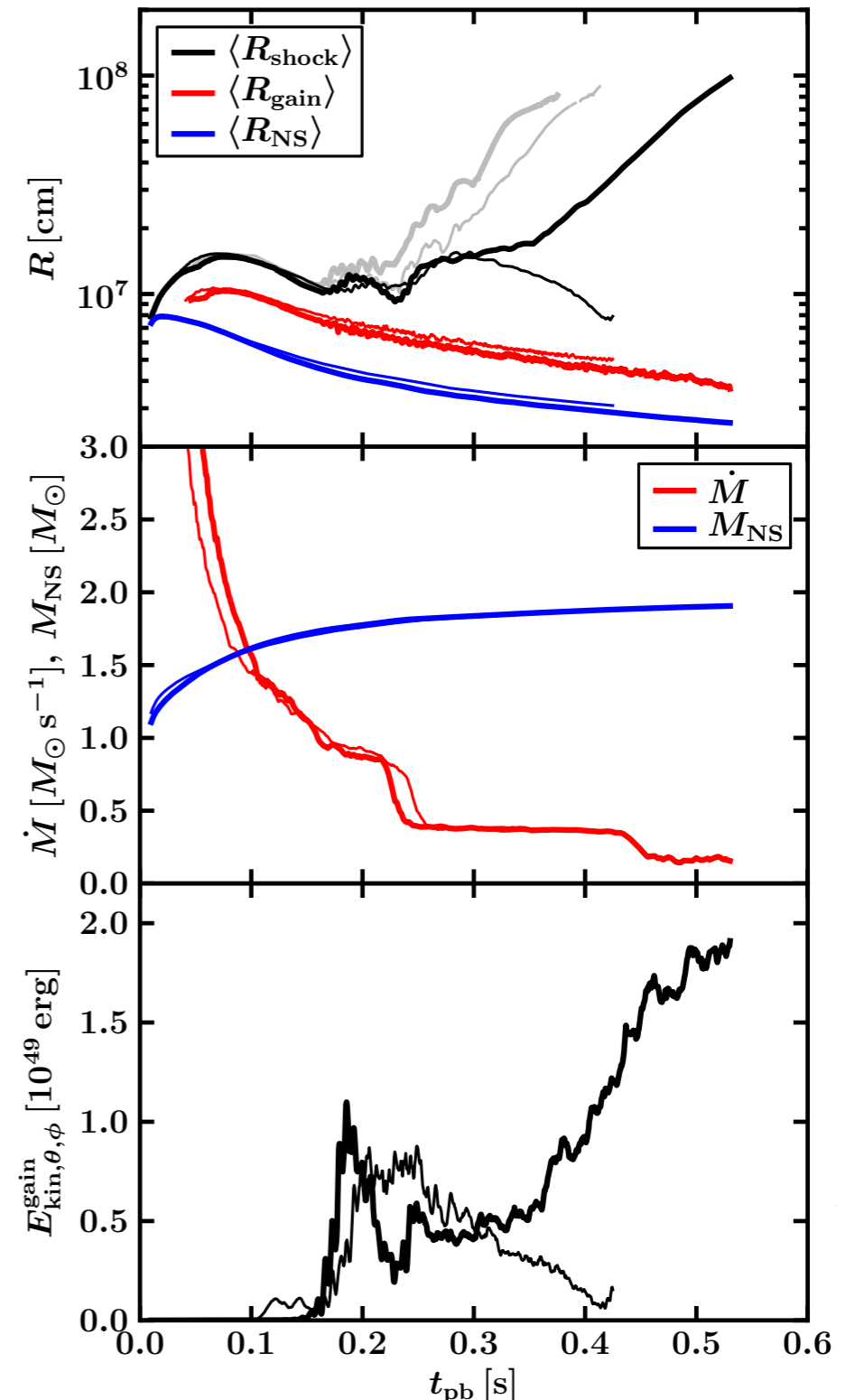
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Melson et al. '15

Two Moment Neutrino Transport

Boltzmann Equation:

$$\frac{\partial x^\alpha}{\partial \tau} \frac{\partial f(x^\mu, p^\mu)}{\partial x_\alpha} + \frac{\partial p^i}{\partial \tau} \frac{\partial f(x^\mu, p^\mu)}{\partial p_i} = \tilde{S}(x^\mu, p^\mu)$$

Take angular moments of the neutrino distribution function:


$$M_{(v)}^{A_k} = \int dV_p \frac{p^{\alpha_1} \dots p^{\alpha_k}}{(-p_\mu u^\mu)^{k-2}} f(p^\beta, x^\beta) \delta(v + p_\delta u^\delta)$$

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Get conservation equations for projections of the rest frame energy dependent stress tensor:

$$M_{(\nu)}^{\alpha\beta};\beta$$



$$\begin{aligned} \partial_t \tilde{E} + \partial_j (\alpha \tilde{F}^j - \beta^j \tilde{E}) + \partial_\nu (\nu \alpha n_\alpha \tilde{M}^{\alpha\beta\gamma} u_{\gamma;\beta}) &= \alpha [\tilde{P}^{ij} K_{ij} - \tilde{F}^j \partial_j \ln \alpha - \tilde{S}^\alpha n_\alpha] \\ \partial_t \tilde{F}_i + \partial_j (\alpha \tilde{P}_i^j - \beta^j \tilde{F}_i) - \partial_\nu (\nu \alpha \gamma_{i\alpha} \tilde{M}^{\alpha\beta\gamma} u_{\gamma;\beta}) &= \alpha \left[\frac{\tilde{F}_k \partial_i \beta^k}{\alpha} - \tilde{E} \partial_i \ln \alpha + \frac{\tilde{P}^{jk}}{2} \partial_i \gamma_{jk} + \tilde{S}^\alpha \gamma_{i\alpha} \right] \end{aligned}$$

Two Moment Neutrino Transport

Boltzmann Equation:

$$\frac{\partial x^\alpha}{\partial \tau} \frac{\partial f(x^\mu, p^\mu)}{\partial x_\alpha} + \frac{\partial p^i}{\partial \tau} \frac{\partial f(x^\mu, p^\mu)}{\partial p_i} = \tilde{S}(x^\mu, p^\mu)$$

Take angular moments of the neutrino distribution function:

$$\rightarrow M_{(\nu)}^{A_k} = \int dV_p \frac{p^{\alpha_1} \dots p^{\alpha_k}}{(-p_\mu u^\mu)^{k-2}} f(p^\beta, x^\beta) \delta(\nu + p_\delta u^\delta)$$

Get conservation equations for projections of the rest frame energy dependent stress tensor:

$$M_{(\nu)}^{\alpha\beta} ; \beta \rightarrow \begin{aligned} \partial_t \tilde{E} + \partial_j (\alpha \tilde{F}^j - \beta^j \tilde{E}) + \partial_\nu (\nu \alpha n_\alpha \tilde{M}^{\alpha\beta\gamma} u_{\gamma;\beta}) &= \alpha [\tilde{P}^{ij} K_{ij} - \tilde{F}^j \partial_j \ln \alpha - \tilde{S}^\alpha n_\alpha] \\ \partial_t \tilde{F}_i + \partial_j (\alpha \tilde{P}_i^j - \beta^j \tilde{F}_i) - \partial_\nu (\nu \alpha \gamma_{i\alpha} \tilde{M}^{\alpha\beta\gamma} u_{\gamma;\beta}) &= \alpha \left[\frac{\tilde{F}_k \partial_i \beta^k}{\alpha} - \tilde{E} \partial_i \ln \alpha + \frac{\tilde{P}^{jk}}{2} \partial_i \gamma_{jk} + \tilde{S}^\alpha \gamma_{i\alpha} \right] \end{aligned}$$

Amenable to finite volume techniques

Still need to specify neutrino stress tensor:

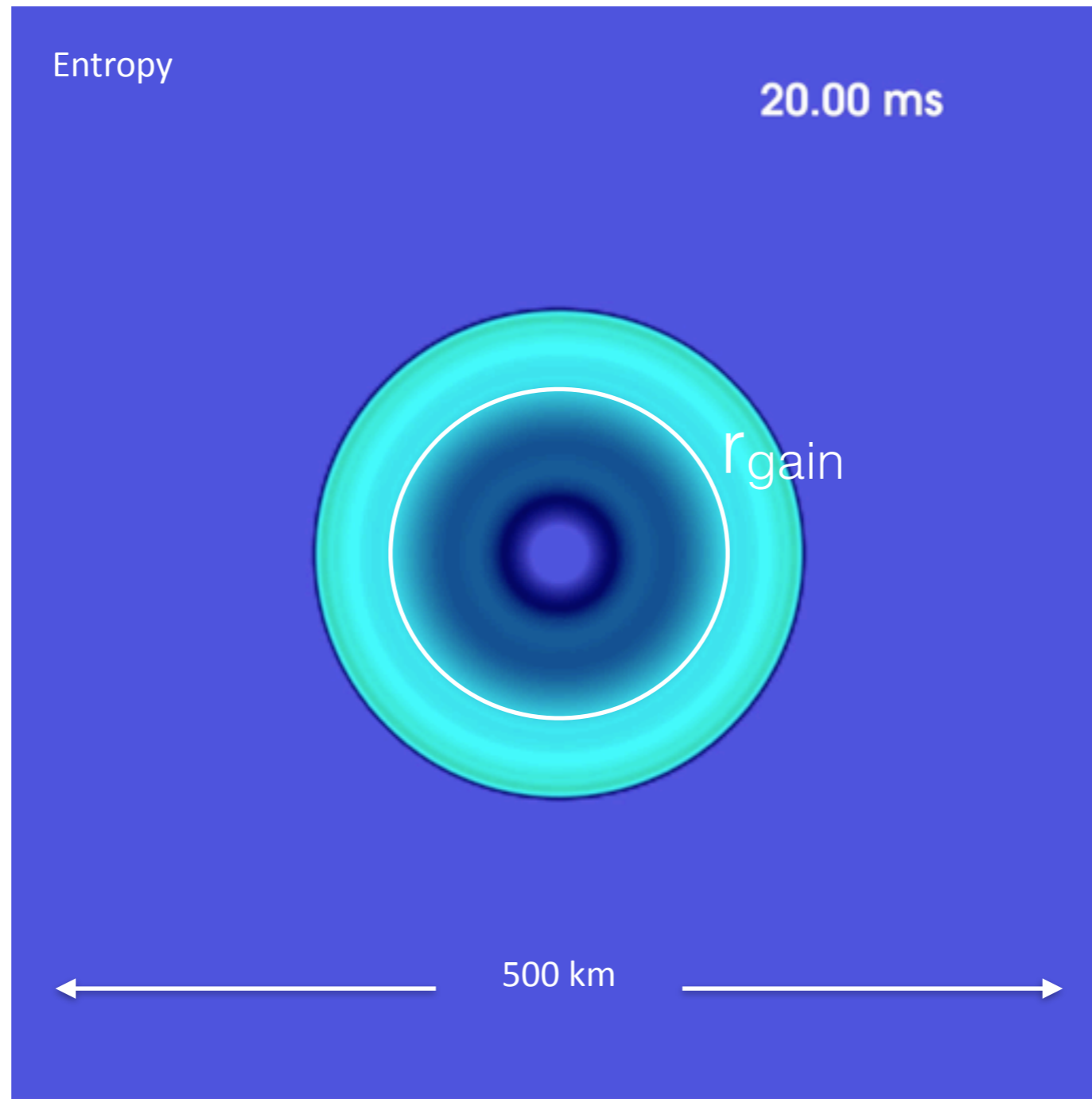
$$P_{(\nu)}^{\alpha\beta} = \frac{3\chi(\xi) - 1}{2} P_{(\nu),thin}^{\alpha\beta} + \frac{3(1 - \chi(\xi))}{2} P_{(\nu),thick}^{\alpha\beta}$$



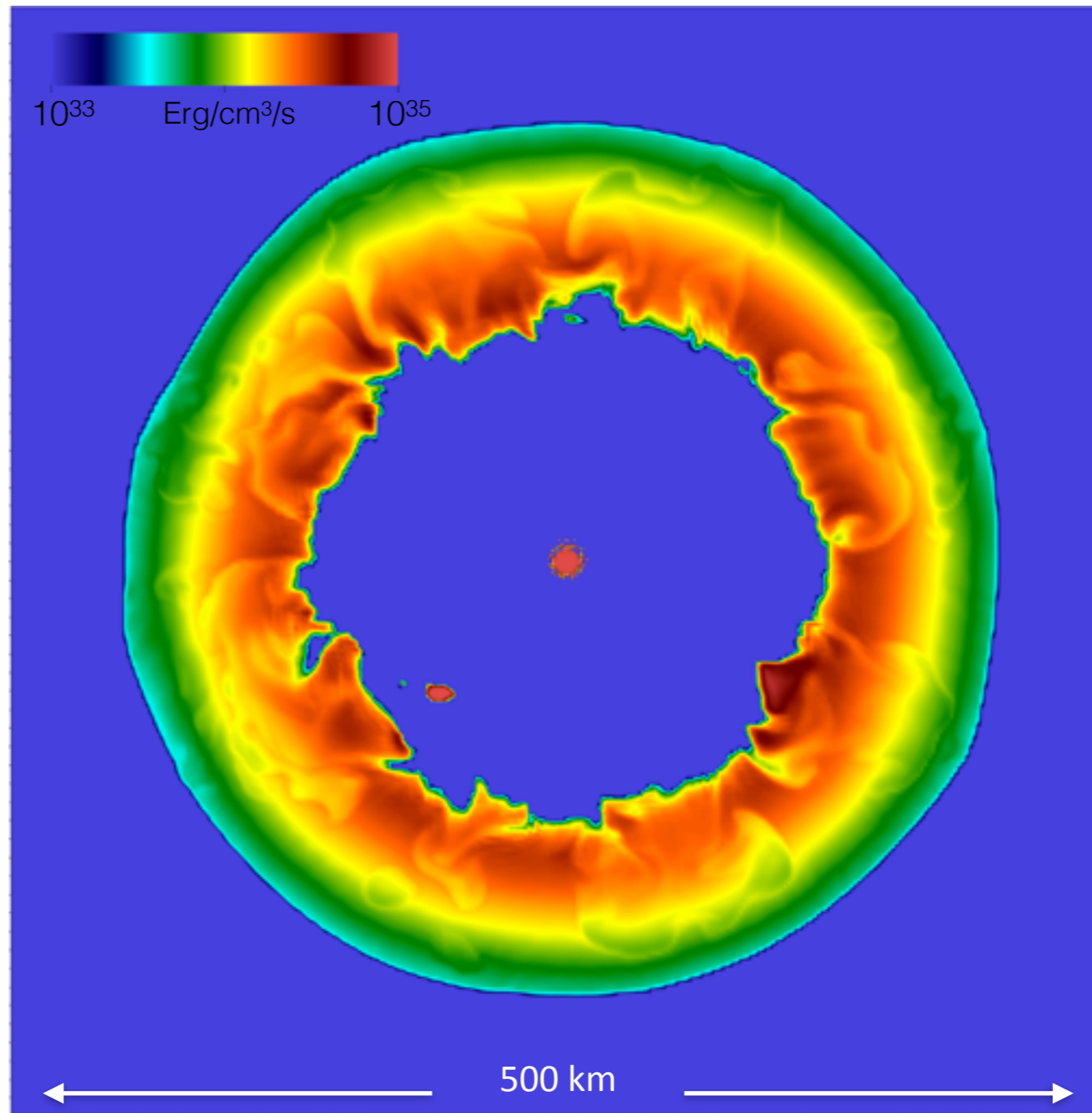
Neutrino Interactions

Reaction	Details & References
$\nu + N \leftrightarrow \nu + N$	Inelastic, non-interacting (Bruenn '85)
$\nu + e^- \leftrightarrow \nu + e^-$	Ultra-relativistic, elastic (Yueh & Buchler '77, etc.)
$\nu + \bar{\nu} + N + N \leftrightarrow N + N$	Non-relativistic, One-pion exchange, uncertain (Hannestad & Raffelt '98)
$\nu + \bar{\nu} \leftrightarrow e^- + e^+$	Ultra-relativistic (Bruenn '85)
$\nu_e + n \leftrightarrow e^- + p$	Inelastic, non-interacting, degeneracy (Bruenn '85)
$\bar{\nu}_e + p \leftrightarrow e^+ + n$	Inelastic, non-interacting, degeneracy (Bruenn '85)

Post Bounce Evolution of CCSNe



Neutrino Heating Rate



Conclusions

- M1 provides a reasonably good method for radiative transfer in quasi-spherical situations
- Long term evolution of post-bounce CCSNe evolution without any imposed symmetries
- Shock runaway occurs in many models, but have not quantified predicted explosion energies
- Significant dependence on resolution and assumed symmetries
- More detailed analysis of post-shock flow required

